

The Karhunen-Loeve Orthagonal Expansion

Suppose $\phi_n(t)$ are a set of orthonormal basis in the interval (0,T). A function X(t) (deterministic or random) may be expanded as

$$X(t) = \sum c_n \varphi_n(t), \qquad 0 < t < T, \qquad (1)$$

where the coefficient c_n are given by

$$c_n = \int_0^T X(t) \varphi_n(t) dt .$$
⁽²⁾

Note that the property

$$\int_{0}^{T} \varphi_{n}(t) \varphi_{m}^{*}(t) dt = \delta_{nm}$$
(3)

was used in the derivation of (2).

When X(t) is a random function, the coefficients c_n become random coefficients) variables. In the following, assume $E\{X\}=0$.

Theorem: In the expansion (1), the coefficients c_n become uncorrelated (orthogonal) random variables if and only if $\varphi_n(t)$ are the eigenfunction of the following Fredholm's integral equation:

$$\int_0^T \mathbf{R}_{xx}(\mathbf{t}_1, \mathbf{t}_2) \boldsymbol{\varphi}_n(\mathbf{t}_2) = \lambda_n \boldsymbol{\varphi}_n(\mathbf{t}_1).$$
(4)

In this case,

$$\mathbf{E}\left\{\mathbf{c}_{n}\right|^{2}\right\} = \lambda_{n} \,. \tag{5}$$

Proof: From (1) and (2), it follows that

$$E\{x^{*}(t_{1})c_{n}\} = E\{c_{n}|^{2}\}\phi_{n} = \int_{0}^{T}R_{xx}(t_{1},t_{2})\phi_{n}(t_{2})dt_{2}, \qquad (6)$$

where

$$\mathbf{E}\left\{\mathbf{c}_{n}\mathbf{c}_{m}^{*}\right\} = \mathbf{E}\left\{\left|\mathbf{c}_{n}\right|^{2}\right\} \delta_{nm}$$

$$\tag{7}$$

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is used. Thus $\lambda_n = E\left\{c_n\right\}^2$ in equation (4).

It may be also shown that the Karhunen-Loeve (K-L) expansion converges in mean-square sense, i.e.

$$\mathbf{E}\left\{\left[X(t) - \sum_{n} c_{n} \boldsymbol{\varphi}_{n}(t)\right]^{2}\right\} = 0.$$
(8)

(See Papoulis page 304 for details.) It may also be easily shown that

$$\mathbf{R}_{xx}(\mathbf{t}_{1},\mathbf{t}_{2}) = \sum_{n} \lambda_{n} |\boldsymbol{\varphi}_{n}|^{2} .$$
⁽⁹⁾

Stationary and Periodic Processes

If X(t) is stationary, then $R_{xx} = R_{xx}(t_1 - t_2)$. If in addition, X(t) is also periodic in the mean-square sense, then

$$\varphi_{n}(t) = \frac{1}{\sqrt{T}} e^{in\omega_{0}t}, \qquad \omega_{0} = \frac{2\pi}{T}.$$
(10)

The K-L expansion for x(t) and R_{xx} are given as

$$\mathbf{x}(t) = \sum_{-\infty}^{+\infty} \frac{c_n}{\sqrt{T}} e^{in\omega_0 t}, \quad \mathbf{E}\left\{c_n\right\}^2 = \lambda_n$$
(11)

$$R_{xx}(t_{1},t_{2}) = \frac{1}{T} \sum_{-\infty}^{+\infty} \lambda_{n} e^{in\omega_{0}(t_{1}-t_{2})}$$
(12)

The power spectrum of X(t) is then given by

$$S_{xx}(\omega) = \frac{1}{T} \sum_{-\infty}^{+\infty} \lambda_n \delta(\omega - n\omega_0).$$
⁽¹³⁾

Furthermore,

$$E\{X^{2}(t)\} = \frac{1}{T} \sum_{-\infty}^{+\infty} \lambda_{n} . \qquad (14)$$

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Stationary Nonperiodic Processes

These may be considered to have infinite period. One may write

$$X(t) = \int_{-\infty}^{+\infty} e^{i\omega t} n(\omega) \sqrt{S(\omega)} d\omega, \qquad (15)$$

where $n(\omega)$ is a white noise in frequency space with

$$\mathbf{E}\{\mathbf{n}(\boldsymbol{\omega}_1)\mathbf{n}(\boldsymbol{\omega}_2)\} = \delta(\boldsymbol{\omega}_1 - \boldsymbol{\omega}_2) \tag{16}$$

Autocorrelation of X(t) then is given as

$$R_{xx}(t_1,t_2) = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} e^{i\omega_1 t_1 - i\omega_2 t_2} \delta(\omega_1 - \omega_2) \sqrt{s(\omega_1)s(\omega_2)} d\omega_1 d\omega_2$$

Hence,

$$\mathbf{R}_{xx}\left(\mathbf{t}_{1}-\mathbf{t}_{2}\right)=\int_{-\infty}^{+\infty}e^{-i\omega(\mathbf{t}_{2}-\mathbf{t}_{1})}\mathbf{s}(\boldsymbol{\omega})\mathbf{d}\boldsymbol{\omega},$$
(17)

as is expected.

Responses of a Linear System to White Excitations

Consider a linear system

$$L_t X(t) = n(t), X(0) = X'(0) = ... = 0,$$
 (18)

with

$$\mathbf{R}_{nn}(\mathbf{t}_{1},\mathbf{t}_{2}) = 2\pi \mathbf{S}_{0}\delta(\mathbf{t}_{1}-\mathbf{t}_{2}).$$
⁽¹⁹⁾

Solution to (18) is given by

$$X(t) = \int_0^t h(t - \tau) n(\tau) d\tau, \qquad (20)$$

where h(t) is the impulse response. i.e.,

$$L_t h(t) = \delta(t).$$
(21)

Multiplying (20) by $X(t_2)$ and apply L_t , after averaging one finds

$$L_{t}R_{xx}(t,t_{2}) = L_{t}\int_{0}^{t}h(t-\tau)E\{n(\tau)X(t_{2})\}d\tau$$

or

$$L_{t}R_{xx}(t,t_{2}) = E\{n(t)X(t_{2})\} = 2\pi S_{0}'h(t_{2}-t).$$
(22)

Restating equation (4) as

$$\int_{0}^{T} \mathbf{R}_{xx}(\mathbf{t},\mathbf{t}_{2})\boldsymbol{\varphi}(\mathbf{t}_{2})d\mathbf{t}_{2} = \lambda\boldsymbol{\varphi}(\mathbf{t})$$
(23)

and applying L_t and using (22), it follows that

$$\int_0^T 2\pi S_0 h(t_2 - t) \varphi(t_2) dt_2 = \lambda L_t \varphi(t)$$
(24)

Operating on (24) with L_{-t} (same as L_t with t being replaced by -t) and using (21), it follows that

$$\lambda L_{-t} L_{t} \phi(t) = \int_{0}^{T} 2\pi S_{0} \delta(t_{2} - t) \phi(t_{2}) dt_{2} = 2\pi S_{0} \phi(t)$$
(25)

Equation (25) is a differential equation for evaluating the eigenfunction $\phi_n(t)$ and eigenvalues λ_n . It may be shown that if L_t involves derivations of the order N, then the following boundary conditions may be used:

$$\varphi^{(i)}(0) = 0$$
 for $i = 0, 1, ..., N-1$ (26)
 $I_{i} = \sigma^{(i)}(1) = 0$ for $i = 0, 1, ..., N-1$ (27)

$$L_t \varphi^{(i)}(t)|_{t=T} = 0$$
 for $i = 0, 1, ..., N-1$ (27)